

$$|\psi\rangle, A, A \quad A|a\rangle = a|a\rangle$$

$$|\langle a|\psi\rangle|^2 = P(a)$$

$|\psi_n\rangle$ Choose P_n

$$P_1 |\langle a|\psi_1\rangle|^2 + P_2 |\langle a|\psi_2\rangle|^2 + \dots$$

$$\langle A \rangle = \sum_n P_n \langle \psi_n | A | \psi_n \rangle$$

$$\text{Tr} = \sum_n \langle \phi_n | B | \phi_n \rangle$$

$$B = \begin{matrix} & |\psi_1\rangle & \langle \psi_2| \\ \hline |\psi_1\rangle & & \\ \langle \psi_2| & & \end{matrix}$$
$$B|\psi\rangle = |\psi_1\rangle \langle \psi_2|\psi\rangle$$

$|\phi_n\rangle$
complete
or w set
of states.

$$\text{Tr } B = \langle \psi_2 | \psi_1 \rangle$$

$$\sum_n \langle \phi_n | \psi \rangle | \phi_n \rangle = | \psi \rangle$$

$$\rho = \sum_n p_n | \psi_n \rangle \langle \psi_n |$$

$$\text{Tr}(\rho A) = \sum_n p_n \langle \psi_n | A | \psi_n \rangle$$
$$= \langle A \rangle$$

$$\text{Tr}(AB) = \text{Tr}(BA)$$

$$\text{Tr}[X, P] = 0 \quad = i \text{Tr} \dot{1} = 0$$

$$T_n = \sum_m \langle \phi_m | A | \phi_m \rangle$$

$$\Rightarrow \sum_m \langle \tilde{\phi}_m | A | \tilde{\phi}_m \rangle$$

$$| \tilde{\phi}_m \rangle = U | \phi_m \rangle$$

$$= \sum_m \langle \phi_m | U^\dagger A U | \phi_m \rangle$$

$$= T_n \quad U^\dagger A U = T_n A \underbrace{U^\dagger U}_I$$

ρ is playing role of the state
in expectation values.

System of two parts.
Look at operators only on
one of parts
State of whole

$$|\psi\rangle = \sum_n |\psi_n\rangle_1 |\tilde{\psi}_n\rangle_2$$

$$A_1 \langle A \rangle = \sum_n \langle \tilde{\psi}_n | \langle \psi_n | A | \psi_n \rangle | \tilde{\psi}_n \rangle$$
$$= \sum_n \langle \tilde{\psi}_n | \tilde{\psi}_n \rangle \langle \psi_n | A | \psi_n \rangle$$

$$\rho = \sum_m \langle \hat{\Psi}_m | \hat{\Psi}_m \rangle |\Psi_m\rangle \langle \Psi_m|$$

Hermitian $\rho^\dagger = \rho$.

Reduced Density matrix

$$T_V(\rho) = \sum_m \langle \hat{\Psi}_m | \hat{\Psi}_m \rangle \langle \Psi_m | \Psi_m \rangle$$

$$\langle \Psi | \Psi \rangle = 1.$$

Eigenvalues are all positive.

$|\lambda\rangle, \lambda \rightarrow$ eigenstate of ρ

$$\Rightarrow \sum_k \lambda_k = 1, \quad \lambda_k \text{ are } k \text{ probabilities}$$

The reduced density matrix
is exactly like the
"prob. choice" of states.

So a pure state for system
is like a prob. choice for
the subsystem.

Any density matrix, ρ_{sub} , can
model it by a pure state for
that system plus another
system.

Let's say we have system has a pure state

$$\rho = \sum_n \frac{\langle \tilde{\Psi}_n | \Psi \rangle \langle \Psi | \tilde{\Psi}_n \rangle}{\langle \tilde{\Psi}_n | \tilde{\Psi}_n \rangle}$$

If we choose $|\tilde{\Psi}_n\rangle$ to be the eigenstates of ρ

$$\rho = \sum_n \lambda_n \frac{|\tilde{\Psi}_n\rangle \langle \tilde{\Psi}_n|}{\langle \tilde{\Psi}_n | \tilde{\Psi}_n \rangle}$$

$$\langle \tilde{\Psi}_m | \tilde{\Psi}_n \rangle = \delta_{mn} \lambda_n$$

Thus the expansion factors for system 2 must also be an orthonormal system of states

and density matrix

$$\rho_2 = \sum \lambda_n |\tilde{\psi}_n\rangle \langle \tilde{\psi}_n|$$

→ eigenvalues of ρ_2 are same as eigenvalues of ρ_1

→ Rank of $\rho_1 = \#$ values which are non zero

= Rank of ρ_2

If Rank $\rho_1 = 1$
then state is product
state

$$|\psi\rangle = |\psi_1\rangle_1 |\hat{\psi}_1\rangle_2$$

$$\langle A \otimes B \rangle = \langle \psi_1 | A | \psi_1 \rangle \langle \hat{\psi}_1 | B | \hat{\psi}_1 \rangle$$

No correlations.

Rank of $\rho \Rightarrow$ measure of
correlations.

Popular measure:
Entropy.

Start with

$$S = \sum_i -p_i \ln p_i$$

Von Neumann entropy \Rightarrow ρ density matrix

$$\sum_i -\lambda_i \ln \lambda_i$$

for density matrix

$$\text{If } p_i = \frac{1}{D}$$

where D is dimension of Hilbert space, $S = \sum_{i=1}^D -\frac{1}{D} \ln \frac{1}{D} = \ln D$

Note that the statement
 that the e-values of
 ρ_1, ρ_2 are same is true
 only for pure states

For example, if

$$\rho = \frac{1}{2} (|1\rangle_1 |1\rangle_2 \langle -1|_1 \langle 1|_2 + |1\rangle_1 |1\rangle_2 \langle +1|_1 \langle 1|_2)$$

$$\rho_1 = \frac{1}{2} (|1\rangle \langle 1|)$$

$$\rho_2 = \frac{1}{2} (|1\rangle \langle -1| + \frac{1}{2} |1\rangle \langle 1|)$$

Decoherence

Pure state for bipartite system
2 subsystems.

Product state

$$|\psi\rangle_1 \underbrace{|\tilde{\psi}\rangle_2}$$

2 level system. $|\psi_1\rangle$ were a state
of σ_3 , $|1\rangle \Rightarrow \sigma_1 \Rightarrow 50-50 \text{ prob}$
of getting ± 1 or -1 .
bu we would get 100% prob
of getting ± 1 for σ_3 .

If started with $\frac{|1\rangle + |-1\rangle}{\sqrt{2}}$

then $50-50$ for σ_3, σ_2 , and
 $100\% - 0$ for σ_1

Interferer.



Interaction which creates an
entangled state

$$H_2 \rightarrow \left(\frac{|1\rangle_1 + |-1\rangle_1}{\sqrt{2}} \right) |1\rangle_2$$

\Downarrow entanglement

$$\left(\frac{|1\rangle_1 |1\rangle_2 + |-1\rangle_1 |-1\rangle_2}{\sqrt{2}} \right) \sigma_1 \sigma_2$$

Density matrix for system 1

$$\rho_1 = \frac{1}{2} \begin{pmatrix} |1\rangle_1 \langle 1|_1 + \frac{1}{2} |-1\rangle_1 \langle -1|_1 \end{pmatrix}$$

For $\frac{1}{2} I$ the prob of getting
 ± 1 for any of $\sigma_1, \sigma_2, \sigma_3$

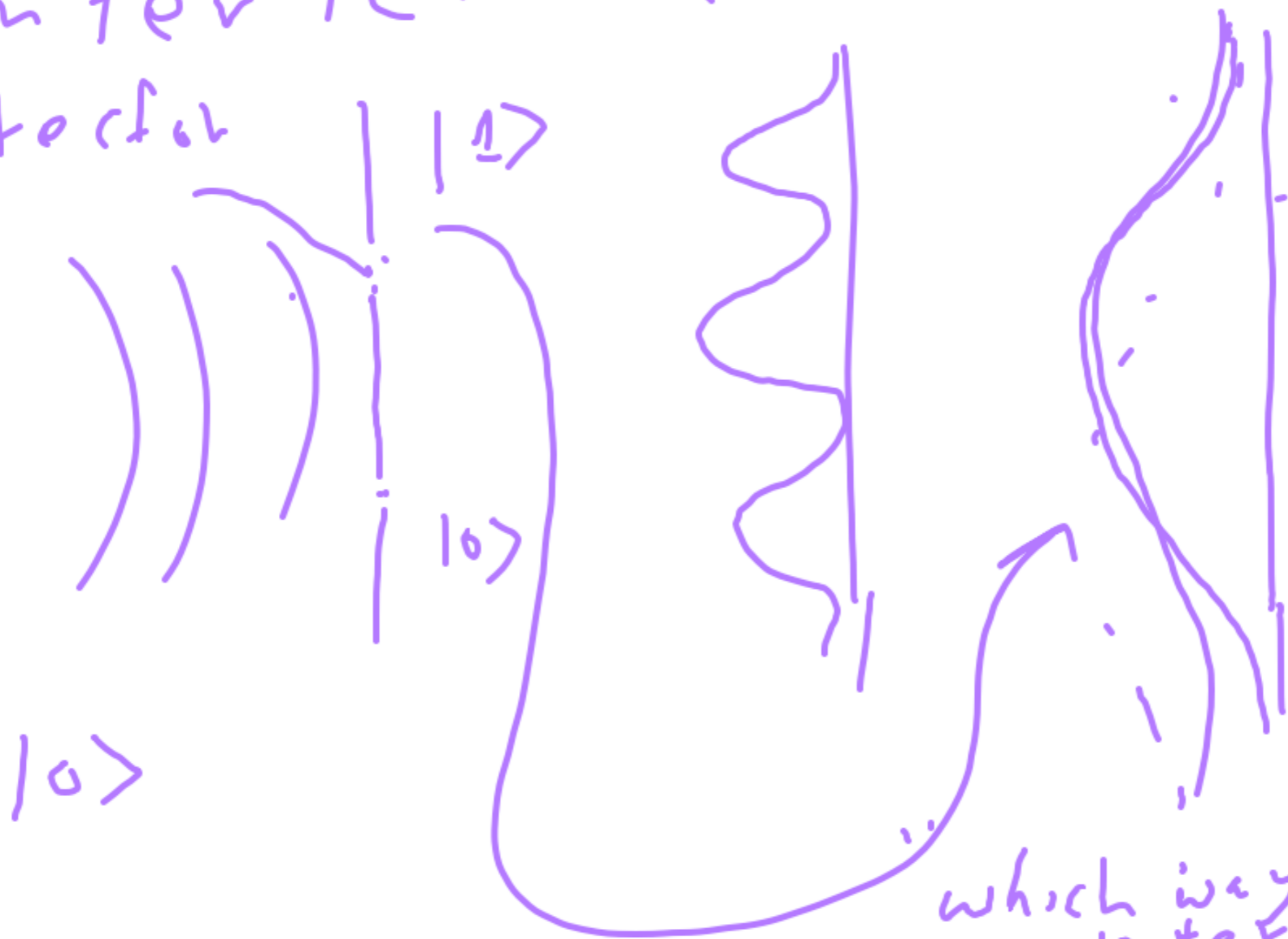
$$= \frac{1}{2}$$

$$\begin{aligned} I &= \rightarrow |1\rangle\langle 1| + |-1\rangle\langle -1| \\ &= \left(\frac{1}{\sqrt{2}} (|1\rangle + |-1\rangle) \right) \left(\langle 1| + \langle -1| \right) \\ &\quad + \left(\frac{1}{\sqrt{2}} (|1\rangle - |-1\rangle) \right) \left(\langle 1| - \langle -1| \right) \end{aligned}$$

\Rightarrow No interference.

Correlations destroy
interference.

Detector



which way " destroys
interference pattern

Quantum Eraser