

0.1 Kruscal Schwartzschild

The Schwartzschild metric is

$$ds^2 = \left(1 - \frac{2M}{r}\right)dvdu - r^2d\Omega^2 \quad (1)$$

where $d\Omega^2 = d\theta^2 + \sin(\theta)^2d\phi^2$ and

$$u = t - r^* \quad (2)$$

$$v = t + r^* \quad (3)$$

$$r^* = r + 2 * M \ln\left(\frac{r}{2M} - 1\right) \quad (4)$$

Thus $v - u = 2r^* + 4 \ln\left(\frac{r-2M}{2M}\right)$ and $e^{(v-u)/4M} = \left(\frac{r-2M}{2M}\right)e^{r/2M}$ and we can write

$$ds^2 = \frac{2M}{r}e^{-r/2M}e^{(v-u)/4M}dudv - \left(\frac{v-u}{2} - 2M \ln\left(\frac{r}{2M} - 1\right)\right) \quad (5)$$

We now define

$$\begin{aligned} U &= -4Me^{-u/4M} & u > 0 \\ V &= 4Me^{v/4M} & v < 0 \end{aligned} \quad (6)$$

Now define

$$V = \tan(\hat{V}/4M) \quad (7)$$

$$U = \tan(\hat{U}/4M) \quad (8)$$

For large r , $v - u = 2r^* = 2r\left(1 + 2M\frac{\ln(r/2M-1)}{r}\right)$ This approaches r for large r , but the $2M \ln(r/2M-1)$ is a singular function as $r = \frac{\sin(V)\sin(-U)}{\cos(V)\cos(U)}$ goes to infinity. Thus the $U = -\pi/2$ and $V = \pi/2$ surfaces are well behaved but the metric as one approaches them is not very well behaved. and is especially badly behaved as one approaches both of those limits. I_0 , spacelike infinity, is badly behaved not only in the conformal factor going to infinity there, but also that the conformal metric is badly behaved there. Similarly, the metric at I^\pm is badly behaved unlike the case for flat spacetime.

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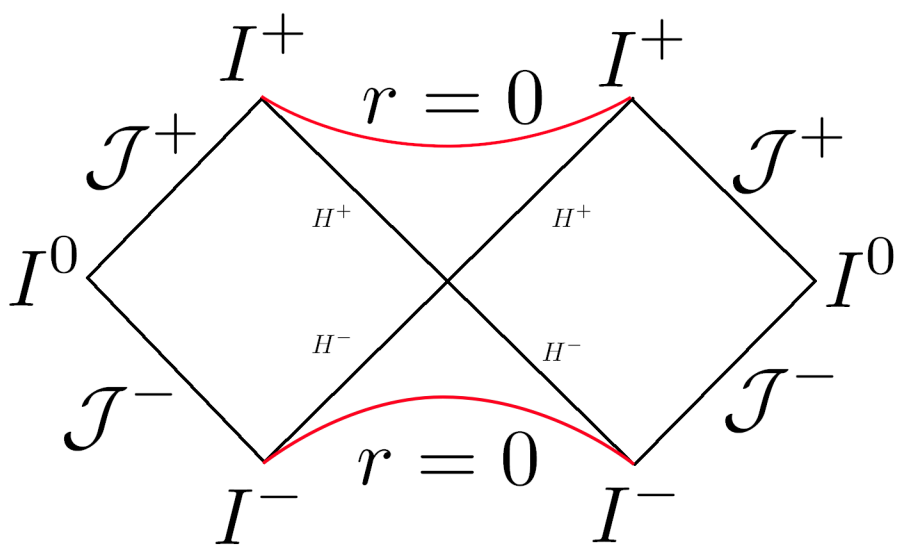


Figure 1: The Penrose conformal diagram for the eternal (Kruskal) black hole.